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# Late Pleistocene evolution of tides and tidal dissipation

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## Key Points:

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| 7  | • | Tides over the last glacial cycle are explicitly modelled and tidal dissipation ranging back |
|----|---|--|
| 8  |   | to 430 ka is inferred with regression analysis.  |
| 9  | • | Enhanced open ocean dissipation (1.8–2.5 $\times$ present day) occurred during glacial max-  |
| 10 |   | ima, and near-present-day values during interglacials.                                       |
| 11 | • | Peak glacial M2 tidal dynamics are very sensitive to changes in ice sheet extent, which may  |
| 12 |   | influence ocean mixing and glacial climate.  |
|    |   |  |

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#### 13 Abstract

Studies of the Last Glacial Maximum (LGM; 26.5-19 ka) tides showed strong enhancements in 14 open ocean tidal amplitudes and dissipation rates; however, changes prior to the LGM remain largely 15 unexplored. Using two different ice sheet and sea level reconstructions, we explicitly simulate the 16 evolution of the leading semi-diurnal and diurnal tidal constituents  $(M_2, S_2, K_1 \text{ and } O_1)$  over the 17 last glacial cycle with a global tide model. Both sets of simulations show that global changes, 18 dominated by the Atlantic, take place for the semi-diurnal constituents, while changes for the diurnal 19 constituents are mainly regional. Irrespective of the reconstruction, open ocean dissipation peaks 20 during the sea level lowstands of MIS 2 ( $\sim$ 20 ka) and MIS 4 ( $\sim$ 60 ka), although dissipation values 21 prior to MIS 2 are sensitive to differences in reconstructed ice sheet extent. Using the statistically 22 significant relationship between global mean sea level and dissipation, we apply regression analysis 23 to infer open ocean and shelf dissipation, respectively, over the last four glacial cycles back to 430 24 ka. Our analysis shows that open ocean tidal energy was probably increased for most of this period, 25 peaking during glacial maxima, and returning to near-present-day values during interglacials. Due to 26 tidal resonance during glacial phases, small changes in bathymetry could have caused large changes 27 in tidal amplitudes and dissipation, emphasising the need for accurate ice margin reconstructions. 28 During glacial phases, once global mean sea level decreased by more than  $\sim 100$  m, the amount of 29 open ocean tidal energy available for ocean mixing approximately doubled. 30

### 31 **1 Introduction**

Tides are important for numerous processes in the ocean: in coastal areas, they shape intertidal 32 ecosystems and morphology. In shelf-sea areas, tidal dynamics determine the location of tidal mixing 33 fronts, which separate seasonally stratified waters from year-round mixed waters (e.g., Simpson & 34 Pingree, 1978). This partitioning is important for shelf sea ecosystems, biogeochemical cycles and 35 the export of  $CO_2$  from the shelf seas to the deep ocean (Thomas et al., 2004). In the open ocean, 36 tidally driven mixing supplies approximately half ( $\sim 1$  TW) of the energy necessary to sustain the 37 large-scale meridional overturning circulation (Wunsch & Ferrari, 2004; Ferrari & Wunsch, 2009). 38 Recent work (Wilmes et al., 2021) demonstrated that, for the the Last Glacial Maximum (26.5–19 39 thousand years before present (ka); the LGM hereafter), strong increases in tidal mixing (due to the 40 sea level lowstand and associated changes in tidal dynamics) could be constrained from sediment 41 carbon isotopes. As sea level index points (SLIPs) in coastal areas are generally related to a given 42 high or low tide level and not mean sea level, knowledge of past tidal range changes is important 43 for reconstructing past sea levels (see e.g., Ward et al., 2016). Furthermore, marine terminating ice 44 sheet dynamics are affected by tidal dynamics which influence grounding line movement (Milillo 45 et al., 2017; Batchelor et al., 2023), basal melting (Milillo et al., 2017; Anselin et al., 2023), ice 46 shelf flexure (Walker et al., 2013) and ice flow (Bindschadler et al., 2003; Anandakrishnan, 2003; 47 Gudmundsson, 2007). 48

Reconstructions of global tides and tidal dissipation during the Quaternary (2.58 Ma to present) 49 have generally focused on the last  $\sim\!25$  thousand years (kyr) encompassing the LGM, the deglacial 50 (19–11.7 ka), and the Holocene (11.7 ka to present) (Egbert et al., 2004; Uehara et al., 2006; Griffiths 51 & Peltier, 2008, 2009; Green, 2010; Wilmes & Green, 2014; Wilmes et al., 2019, 2022; Sulzbach et 52 al., 2023). These investigations showed surprising results: the tides were strongly enhanced in the 53 Atlantic during the LGM, especially in the semi-diurnal band, with tidal energy dissipation (i.e., the 54 loss of energy of the tide to bed friction and to the internal tide) for the M<sub>2</sub> tide a factor 2–3 larger 55 than at present in the open ocean. Changes in the North Atlantic are thought to have been particularly 56 strong with amplitudes tripling with respect to present and exceeding 6 m in the Labrador Sea during 57 the LGM (Griffiths & Peltier, 2008, 2009; Wilmes & Green, 2014). These amplifications resulted 58 from changes in ocean basin shape driven by the  $\sim 130$  m global mean sea level (GMSL) drop and 59 associated increases in ice sheet extent. Together, these factors rendered the Atlantic more conducive 60 to resonant amplification of the semi-diurnal tides, thus leading to substantial increases in amplitudes 61 and tidal energy dissipation (e.g., Egbert et al., 2004; Green, 2010). Through the deglacial and the 62

Holocene, energy losses in the deep ocean decreased dramatically whilst the shelf seas (which were 63 emersed during the LGM) re-flooded and became more tidally energetic. 64

On longer time scales, covering the middle and late Pleistocene (~770-11.7 ka), sea level 65 fluctuated by 130-145 m as climate cycled between glacial and interglacial phases (e.g., Fox-Kemper 66 et al., 2021). During this period, ice sheet extent followed a saw-tooth shaped pattern: glacial phases 67 were generally characterised by a long-term gradual cooling culminating in a glacial maximum with 68 peak ice sheet extent and sea level lowstands. Subsequently, climate warmed rapidly and transitioned 69 to interglacial conditions, with climate similar to, or warmer than the pre-industrial. GMSL attained 70 highstands of up to +15 m and ice sheet extent was similar to, or less than present. The last glacial 71 period spanned from the end of the Last Interglacial (LIG; 130-115 ka; Marine Isotope Stage (MIS) 72 5e) to the onset of the Holocene, with significant variations in global mean sea level (Waelbroeck et 73 al., 2002; Siddall et al., 2003; Lambeck, 2004; Spratt & Lisiecki, 2016)(see Figure 1a). During the 74 early part of the glacial period (MIS 5 d-a; 115–71 ka), GMSL fluctuated between relative high- and 75 lowstands of -9 and -50 m, respectively (e.g., Creveling et al., 2017). During MIS 4 (71-57 ka), 76 a GMSL lowstand of around -80 m was reached. Thereafter, GMSL rose to a relative highstand 77 during MIS 3 (57–29 ka), though the exact magnitude remains debated (e.g., Dalton et al., 2022). 78 The GMSL lowstand of  $\sim -130$  m during the LGM (MIS 2) was reached between 26.5 and 20 ka 79 (Clark et al., 2009; Lambeck et al., 2014; Peltier et al., 2015; Gowan et al., 2021). At the onset of 80 the deglacial period (19–11.7 ka), GMSL first rose gradually, then more rapidly by around 100 m 81 until the early Holocene ( $\sim$ 8 ka) when present-day levels were reached. Because tides are sensitive 82 to water depth and changes in ocean basin shape and they behave like shallow water waves, it is 83 expected that they are affected by these large sea level changes, as previously seen for the period 84 spanning from the LGM to the present (e.g., Egbert, 2004; Uehara et al., 2006; Green, 2010; Wilmes 85 & Green, 2014; Wilmes et al., 2022; Sulzbach et al., 2023). 86

In this work, we aim to extend our knowledge of Pleistocene tidal dynamics back to  $\sim$ 430 ka. 87 First, we explicitly model tides over the past glacial cycle covering the period from the LIG to present 88 using two different ice sheet and sea level reconstructions, thus expanding the work in Wilmes et al. 89 (2022) to the entire last glacial cycle. Second, because no spatially and temporally highly resolved 90 global sea level and ice sheet reconstructions exist for the previous multiple glacial cycles, we use the 91 last glacial cycle simulations together with uniform sea level change simulations and extrapolate tidal 92 energy dissipation back to 430 ka based on linear regression analysis. We aim to (1) produce global 93 spatially-varying dissipation estimates for the last glacial cycle which can be used for modelling of the 94 late Pleistocene climate and ocean circulation; and (2) improve our understanding of late Pleistocene 95 tidal dynamics which is relevant for, e.g., SLIPs, ice sheet dynamics, or shelf sea oceanographic 96 processes. 97

#### 2 Methods 98

#### 2.1 Tide model

For the tide model simulations, we use the Oregon State Tidal Inversion Software (thereafter 100 OTIS; see Egbert et al., 2004; Green & Nycander, 2013; Wilmes et al., 2019, for details) in its 101 forward mode. OTIS has been used for numerous paleotide (e.g., Egbert et al., 2004; Wilmes & 102 Green, 2014; Green & Huber, 2013; Green et al., 2017) and future tide applications (e.g., Carless 103 et al., 2016; Wilmes et al., 2017; Pickering et al., 2017; Hayden et al., 2020). OTIS solves the 104 linearised shallow water equations (e.g., Hendershott, 1972): 105

$$\frac{\partial \mathbf{U}}{\partial t} + \mathbf{f} \times \mathbf{U} = -gH\nabla(\zeta - \zeta_{EQ} - \zeta_{SAL}) - \mathbf{F}$$
(1)  
$$\frac{\partial \zeta}{\partial t} = -\nabla \cdot \mathbf{U}$$
(2)

(2)

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where U is the depth integrated volume transport, H denotes water depth, f is the Coriolis vector, qis the gravitational constant,  $\zeta$  denotes tidal elevation,  $\zeta_{EQ}$  stands for the equilibrium tidal elevation, and  $\zeta_{SAL}$  is the tidal elevation due to self-attraction and loading (SAL; i.e., the combined effects of gravitational attraction among the water masses, seafloor deformation, and associated changes in the gravity potential).  $\mathbf{F} = \mathbf{F}_B + \mathbf{F}_{IT}$  represents frictional losses due to bed friction ( $\mathbf{F}_B$ ) and tidal conversion ( $\mathbf{F}_{IT}$ ). The former is represented by the standard quadratic law:

$$\mathbf{F}_B = C_d \mathbf{u} |\mathbf{u}| \tag{3}$$

where  $C_d = 0.003$  is a drag coefficient, and u tidal velocity. The energy losses to the internal tide,  $\mathbf{F}_{IT} = C_{IT}\mathbf{U}$ , depend on a conversion coefficient  $C_{IT}$  given by (Zaron & Egbert, 2006; Green & Huber, 2013)

$$C_{IT}(x,y) = \gamma \frac{(\nabla H)^2 N_b N}{8\pi\omega}$$
(4)

where  $\gamma = 37.5$  is a scaling factor (see Zaron & Egbert, 2006, for more details),  $N_b$  is the buoyancy 119 frequency at the sea-bed,  $\bar{N}$  is the vertical average of the buoyancy frequency, and  $\omega$  is the frequency 120 of the tidal constituent under evaluation.  $\gamma$  was tuned following the process described in Wilmes 121 and Green (2014) to both minimise present-day amplitude root-mean square errors against TPXO9 and fit TPXO9 dissipation values. Horizontally uniform abyssal stratification was assumed which is 123 parameterized by the buoyancy frequency N through  $N(z) = N_0 e^{(-z/1300)}$  with  $N_0 = 5.24 \times 10^{-3}$ . 124 Sensitivity simulations in Wilmes and Green (2014) and Schmittner et al. (2015), where the sensitivity 125 to glacial interglacial stratification changes was explored, showed that dissipation is rather insensitive 126 to glacial-interglacial stratification changes. Tidal dissipation associated with the combined action 127 of  $\mathbf{F}_B$  and  $\mathbf{F}_{IT}$  is calculated following the energy balance method outlined in Ray et al. (2003). 128

#### 129 2.2 Simulations

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### 130 2.2.1 Last Glacial Cycle

In our tide model runs, we represent the ice sheet and sea level history over the last glacial 131 cycle with two different ice sheet reconstructions, one external reconstruction by Gowan et al. (2021) and one derived by us, using a global gravitationally self-consistent sea level model. For the latter, 133 we simulate sea level based on the ICE6G C ice history (Argus et al., 2014; Peltier et al., 2015) 134 that spans the period 122 ka to present with a temporal resolution of 250–500 yrs prior to 21 ka, 135 1000 yrs up to 32 ka, and 2000 yrs thereafter (referred to as "ICE-6G" from this point onwards. In 136 the sea level model, we compute the time-varying deformation of a rotating, Maxwell viscoelastic 137 Earth model with a depth-dependent Earth structure (e.g., Kendall et al., 2005; Milne & Mitrovica, 138 1996). For consistency with the global ICE6G\_C ice model, we adopt the corresponding VM5a 139 depth-dependent Earth model to represent global Earth structure (Argus et al., 2014; Peltier et al., 140 2015). Importantly for our application here, the sea level model includes migrating shorelines and 141 the inundation of water into regions previously covered by marine ice (Mitrovica & Milne, 2003). 142 The model formulation is solved using a pseudo-spectral numerical scheme up to spherical harmonic 143 degrees and order 256 (Kendall et al., 2005). 144

As an alternative to the boundary conditions derived from the ICE6G C ice sheet history, 145 we use the ice sheet and sea level reconstruction from Gowan et al. (2021). Gowan et al. (2021) 146 presented a global high-resolution ice sheet reconstruction for the period 80 ka to present. The 147 reconstruction is consistent with ice physics but was inferred independently of far-field sea level 148 and  $\delta^{18}O$  proxy records (i.e., based on ice sheet margins and constructed to be consistent with 149 simple ice sheet physics). The sea level predictions provided by Gowan et al. (2021) also rely on a 150 global gravitationally self-consistent sea level model (SELEN; e.g., Spada & Stocchi, 2007) with a 151 depth-dependent Earth structure, as well as shoreline migration. These reconstuctions have a lower 152 temporal resolution of 2500 yrs. 153

For our study, both ice sheet reconstructions and the associated sea level change fields were interpolated to the finite-difference grid of OTIS (1/8° spacing in both latitude and longitude) and added to the present-day base topography RTopo-2 (Schaffer et al., 2016). We note that both the ice sheet history and the solid Earth structure, and thus sea level, differ between the two approaches. For the ICE-6G sea levels, tide runs were performed at 2000-year intervals. Each simulation was run with M<sub>2</sub>, S<sub>2</sub>, K<sub>1</sub> and O<sub>1</sub> equilibrium tidal forcing and a simplified SAL scheme that sets  $\zeta_{SAL} = \beta \zeta$ , with  $\beta = 0.1$ . The computational grid extends all the way to 89°N, where it is bounded by an artificial

|                | RMSE (cm) |       |        | Modelled dissipation (TW) |       |        | TPXO | TPXO9 dissipation (TW) |        |  |
|----------------|-----------|-------|--------|---------------------------|-------|--------|------|------------------------|--------|--|
|                | deep      | shelf | global | deep                      | shelf | global | deep | shelf                  | global |  |
| $M_2$          | 3.7       | 10.1  | 4.7    | 1.2                       | 1.3   | 2.5    | 0.8  | 1.5                    | 2.4    |  |
| $S_2$          | 2.0       | 4.4   | 2.3    | 0.2                       | 0.3   | 0.5    | 0.2  | 0.3                    | 0.5    |  |
| $\mathbf{K}_1$ | 1.1       | 3.2   | 1.4    | 0.2                       | 0.2   | 0.4    | 0.1  | 0.2                    | 0.3    |  |
| $O_1$          | 1.2       | 2.6   | 1.4    | 0.1                       | 0.1   | 0.2    | 0.1  | 0.1                    | 0.2    |  |

**Table 1.** Present-day control evaluation. Amplitude root-mean square errors (RMSE) against TPXO9 for the deep ocean (h > 500 m), shelf seas (h < 500 m), and the global ocean. Integrated dissipation values (diss.) are also given for the deep ocean (h > 500 m), shelf seas (h < 500 m), and the global ocean for both the present-day control and for TPXO9.

vertical wall eliminating the need for open boundaries. Although these two simplifications (SAL
 and pole cap) potentially suppress resonant behavior of the Arctic Ocean during the LGM (Griffiths
 & Peltier, 2009), the errors incurred in terms of globally integrated dissipation are expected to be
 small (Sulzbach et al., 2023) and simulations with an iterative SAL scheme carried out previously
 (Wilmes et al., 2022) show very simlar results for the deglacial.

Additionally, to separate the effect of GMSL changes from land-ocean boundary changes and effects of non-uniform sea level change, we conducted tidal simulations with uniform sea level changes. Here, sea level was uniformly changed from -135 m to +20 m in steps of 5 m. Again, RTopo-2 at  $1/8^{\circ}$  x  $1/8^{\circ}$  grid spacing was used as the present-day base topography and the simulations were carried out for M<sub>2</sub> and K<sub>1</sub>.

The present-day control simulation (i.e., for 0 ka) was benchmarked against the global tidal solution TPXO9-atlas-v5 (https://www.tpxo.net/global/tpxo9-atlas) by calculating amplitude root-mean square errors and comparing globally integrated dissipation values; see Table 1. Evidently, the model yields realistic solutions for both diurnal and semi-diurnal tidal constituents.

#### 175 2.2.2 Multiple glacial cycles (430 ka–present)

Owing to the lack of global bathymetry, ice sheet and sea level reconstructions spanning multiple 176 glacial cycles, it is currently not possible to explicitly model tidal dynamics prior to the LIG. We 177 therefore take a different, novel approach to infer tidal dissipation prior to the last glacial cycle to gain 178 a first-order understanding of tidal dynamics during this period. We first establish regression models 179 between GMSL and globally integrated deep ocean and shelf sea dissipation rates, respectively, using 180 the three simulation sets described above, i.e., runs for (1) the last glacial cycle based on ICE-6G, 181 (2) the last glacial cycle based on Gowan et al. (2021), and (3) uniform GMSL changes. Using 182 the global sea level reconstruction by Spratt and Lisiecki (2016) (thereafter SL16), each regression 183 model is then used to infer tidal dissipation for open ocean and shelf ocean dissipation back to 430 184 ka. Because GMSL is driven by global ice volume changes, and the ice sheet volume changes over 185 the last multiple glacial cycles generally follow similar patterns (e.g., waxing and waning of the 186 Laurentide and Fennoscandinavian ice sheets) (e.g., Batchelor et al., 2019), we here use GMSL as a 187 proxy for the combined spatially-varying ice volume and sea level changes which drive the changes in 188 the tides. The different bathymetry reconstructions provide a measure of the uncertainty introduced 189 by the differing ice margins. 190

For the regression models, we chose a polynomial regressions of order two between GMSL (sl) and globally integrated tidal dissipation (D) for the deep and shelf ocean  $(D_{deep}$  and  $D_{shelf})$ , respectively, for each tidal constituent. The quadratic fit accounts for non-linear interactions between GMSL and dissipation due to e.g., resonance effects. The relationship between GMSL and tidal

<sup>195</sup> dissipation thus takes the following form:

$$D_{deep} = \beta_{d0} + \beta_{d1}sl + \beta_{d2}sl^2 \tag{5}$$

$$D_{shelf} = \beta_{s0} + \beta_{s1}sl + \beta_{s2}sl^2 \tag{6}$$

$$D_{total} = D_{deep} + D_{shelf} \tag{7}$$

where  $\{\beta_{d0}, \beta_{d1}, \beta_{d2}\}$  and  $\{\beta_{s0}, \beta_{s1}, \beta_{s2}\}$  represent the regression coefficients for deep and shelf 199 dissipation for each set of model simulations, respectively. We calculated regression coefficients 200 for the relationships between GMSL and deep and shelf dissipation for all three sets of tide runs, 201 respectively. If the relationship between  $sl^2$  and D was not significant at the 95% confidence level, 202 the order of the polynomial was reduced to one (i.e., linear regression). The relationship between 203 sl and  $D_{deep}$  and  $D_{shelf}$ , respectively, was calculated separately because, in general, deep and shelf 204 dissipation behaved in an anti-correlated manner in relation to GMSL (i.e., when GMSL decreases, 205 open ocean dissipation increases and shelf dissipation decreases). The globally integrated dissipation  $D_{total}$  is given by the sum of  $D_{deep}$  and  $D_{shelf}$ . 207

Using each regression model and the SL16 GMSL reconstruction, we then calculated time series of late Pleistocene deep, shelf, and total dissipation. Dissipation values were predicted for each constituent ( $M_2$ ,  $S_2$ ,  $K_1$  and  $O_1$ ) as per Eqs. (5)–(7) and associated standard deviations ( $\sigma$ ) were deduced by applying the laws of variance propagation:

$$\sigma_{D_{dur}}^{2} = \sigma_{\beta_{10}}^{2} + \sigma_{\beta_{11}}^{2} sl^{2} + \left(\beta_{d1}^{2} + (2\beta_{d2}sl)^{2}\right)\sigma_{sl}^{2} + \sigma_{\beta_{10}}^{2} sl^{4}$$
(8)

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$$\sigma_{J_{deep}}^2 = -\sigma_{J_{d0}}^2 + \sigma_{J_{d1}}^2 + (\beta_{11}^2 + (\gamma_{d2}^2))^2 s_1^2 + \sigma_{J_{d2}}^2$$

$$\sigma_{D_{shelf}}^2 = \sigma_{\beta_{s0}}^2 + \sigma_{\beta_{s1}}^2 sl^2 + \left(\beta_{s1}^2 + (2\beta_{s2}sl)^2\right)\sigma_{sl}^2 + \sigma_{\beta_{s2}}^2 sl^4 \tag{9}$$

$$\sigma_{D_{total}}^2 = \sigma_{D_{deep}}^2 + \sigma_{D_{shelf}}^2 \tag{10}$$

where *sl* denotes GMSL from the SL16 reconstruction with the provided formal error  $\sigma_{sl}$ , and  $\sigma_{\beta...}$  denotes the standard deviations associated with the respective regression coefficients.

#### 217 3 Results

#### 3.1 The last glacial cycle: LIG to present

#### 3.1.1 Ice sheet and sea level evolution

The two ice sheet and sea level reconstructions show pronounced differences during large parts 220 of the last glacial cycle (Fig. 1a and 2). In the following, we will focus on the differences between 221 the two. For a detailed presentation of the two ice sheet and sea level reconstructions, see Gowan et al. (2021), and Argus et al. (2014), Peltier et al. (2015) and Pedersen et al. (in prep). Whilst a fully 223 developed Laurentide Ice Sheet is present from around 100 ka in ICE-6G, a similarly extensive ice 224 sheet is not formed until after 65 ka in the Gowan reconstruction. This leads to a GMSL offset of 225  $\sim$ 30 m between the two reconstructions during late MIS 5 and early MIS 4 and large differences in 226 shallow shelf sea area. During the MIS 4 peak glaciation, the ICE-6G Laurentide Ice Sheet is slightly 227 more extensive than in Gowan, with the Laurentide Ice Sheet and the Cordilleran Ice Sheet joined up, 228 whereas they remain separated in Gowan. However, the Gowan reconstruction has more extensive 229 ice around Antarctica. Overall, as a result of these differences, ICE-6G sea levels show a slightly 230 lower and earlier MIS 4 GMSL lowstand than Gowan. However, whilst GMSL leading into MIS 4 is 231 similar, there are large regional differences in sea level due to the differences in ice history. The largest 232 differences in ice sheet extent and sea level between the two reconstructions are seen during MIS 233 3. ICE-6G shows extensive Northern Hemisphere ice sheets throughout MIS 3 whereas in Gowan, 234 ice sheet extent is strongly reduced around 40 ka, with the Fennoscandinavian and Cordilleran Ice 235 Sheets mostly melted and the Laurentide Ice Sheet strongly reduced. These discrepancies lead to an 236 offset in GMSL of around 60 m in the middle of MIS 3 and shelf sea area is approximately doubled 237 in Gowan in comparison to ICE-6G. Towards the LGM, both reconstructions show expanding ice 238 sheets and a drop in GMSL. ICE-6G shows slightly more extensive Northern Hemisphere ice sheets, 239 whereas Gowan shows a stronger expansion of the Antarctic ice sheet margins. The ICE-6G sea 240 level lowstand occurs around 26 ka, whereas, for Gowan, it does not take place until 20 ka, and is 241

around 20 m less than for ICE-6G. From 20 ka to present, both simulations show a similar GMSL
 evolution. However, again, there are pronounced differences in regional sea levels (locally, > 50 m
 offsets) due to differences in instantaneous ice sheet loading but also ice history.

#### 3.1.2 Amplitude evolution

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Large changes in tidal amplitudes (regionally in excess of 4 m) occur over the last glacial cycle for both sea level reconstructions, but, owing to the discrepancies in sea level and ice sheet history, there are pronounced offsets between the tidal histories (see Figs. 3 and 4 for total amplitude changes). In the following, we discuss how each constituent contributes to the total amplitude changes, focusing the description on the areas that show the largest changes and laying the emphasis on M<sub>2</sub>, which shows the largest absolute changes.

For the ICE-6G simulations, M<sub>2</sub> amplitudes (Figure S1 in Supporting Information S1) show 252 levels similar to present-day during the LIG. But, M2 amplitudes in the North Atlantic, in the Weddell 253 Sea and to a lesser extent in the South Atlantic and in the Gulf of Panama rapidly increase at the 254 termination of the LIG and onwards. A further increase in amplitudes in these areas occurs from 255 68 ka onward towards the sea level lowstand around 60 ka, where, additionally, the Arctic shows 256 large tidal amplitudes (> 3 m in the Chukchi Sea). After 60 ka, the enhancements persist, albeit at a 257 slightly reduced level. During the transition to the LGM, North Atlantic amplitudes increase further 258 to peak round 20 ka, however, Weddell Sea amplitudes are much reduced in comparison to MIS 3 259 because of the more extensive grounded ice in the bay. During the deglacial, pronounced drops in 260 North Atlantic and Arctic  $M_2$  amplitudes take place, and by 8 ka tidal amplitude reach near-present 261 day levels. 262

For the Gowan simulations, M<sub>2</sub> amplitudes (Figure S5 in Supporting Information S1) between 263 80 ka and 72.5 ka show similar levels to present, apart from considerably larger amplitudes in the 264 Weddell and Ross Sea area. As the sea level lowstand around 60 ka is approached, North Atlantic 265 amplitudes strongly increase. These enhancements persist (albeit at a smaller magnitude) until  $\sim$ 55 266 ka, after which the North Atlantic oscillates between periods of larger and smaller M<sub>2</sub> amplitudes with peaks around 45 ka, 37.5 ka and 32.5 ka. Notable are also strongly increased tides in the Ross 268 Sea between 50 and 40 ka. Large North Atlantic amplitudes develop towards the LGM and during 269 the early to mid deglacial (peaking around 17.5 ka and persisting until 12.5 ka). During LGM, strong 270 enhancements in Arctic tides can also be seen. 271

For the ICE-6G runs,  $S_2$  amplitudes (Figure S2 in Supporting Information S1) are similar to present at the end of the LIG. During the early glacial, Labrador Sea and North Atlantic amplitudes increase, but the increases are less pronounced than for  $M_2$ . Furthermore, large  $S_2$  amplitudes occur in the Coral Sea (NE Australia). The two sea level lowstands (~60 ka and LGM) are characterised by large North Atlantic tides and strong enhancements in Arctic  $S_2$  amplitudes.  $S_2$  amplitudes in the Gowan simulations (Figure S6 in Supporting Information S1) follow a similar picture as  $M_2$ amplitudes albeit with a lower magnitude. Notable is that Arctic  $S_2$  amplitudes enhancements are much reduced in the Gowan simulations and are only present during the middle of the LGM.

In the ICE-6G runs,  $K_1$  and  $O_1$  tidal changes (Figure S4 in Supporting Information S1) are mainly regional. During the sea level lowstands around 62 ka and during the LGM, Pacific shelf seas (Sea of Okhotsk, South China Sea and Banda Sea) become resonant. During the remainder of the last glacial cycle, amplitudes remain close to their present-day levels.

Similarly, both  $K_1$  and  $O_1$  amplitudes are characterised by mainly regional changes in the Gowan simulations (Figures S7 and S8 in Supporting Information S1). During periods of sea level lowstands (~62.5–55 ka and 27.5–15 ka), increased amplitudes can be see around Antarctica and in the shelf seas of the Pacific, whereas other parts of the ocean seem rather insensitive to the large sea level and ocean basin shape changes.

## 289 3.1.3 Dissipation evolution

For  $M_2$ , global open ocean dissipation in the ICE-6G simulation is at near-present-day values 290 during the LIG (Figure 1b). It approximately doubles over a period of  $\sim 10$  kyr between 120 and 291 110 ka to  $\sim$ 2 TW when N and S Atlantic dissipation increases (see Fig. 5). It remains elevated 292 at this level until around 70 ka when dissipation rises by a further 25% (0.5 TW) driven by Arctic 293 and N and S Atlantic dissipation increases. Open ocean dissipation remains at these elevated values 294 peaking around 65 ka, 40 ka and during the LGM, with only a small dip occurring around 28 ka. 295 From 68 ka to 30 ka, open ocean dissipation in the Southern Ocean more than doubled in the ICE-6G runs driven by increases in the Weddell Sea area which increases Southern Hemisphere dissipation to levels greater than in the Northern Hemisphere (Figure S9 in Supplementary Information S1). 298 During MIS 4 to MIS 2, dissipation increases around the European Shelf, in the Labrador Sea, 299 Denmark Strait, Norwegian Sea, along the mid-Atlantic ridge, in the Arctic Basin and also in the 300 South Atlantic. In contrast, the Gowan simulations show near present-day open ocean dissipation 301 values (apart from around Antarctica which increases Southern Hemisphere dissipation; Figure S9 302 in Supplementary Information S1) prior to the sea level lowstand around 60 ka when they increase 303 by  $\sim 50\%$  during the sea level lowstand through Atlantic enhancements (see Fig. 6). Thereafter, dissipation decreases slightly, but oscillates through three distinct troughs and peaks. It is notable 305 that Hudson Bay & Strait dissipation is significantly anti-correlated with North Atlantic dissipation, 306 suggesting that when Hudson Bay & Strait lies dry or is ice covered, North Atlantic dissipation peaks. 307 Towards the LGM, dissipation increases again, and reaches its peak during the early deglacial phase 308 with a 90% increase relative to present driven by increases in Atlantic (especially North Atlantic) 309 and Arctic dissipation, whereas in the ICE-6G simulations, dissipation in the Southern Hemisphere 310 decreases-in contrast to the Gowan simulations. 311

Despite the large differences in shelf area (see Fig. 1a), the evolution of shelf sea dissipation is 312 relatively similar between the two sets of simulations. Between 120 and 70 ka, shelf dissipation is 313 at levels similar to present. With the decrease in sea level towards the lowstand around 60 ka, shelf 314 dissipation decreases by  $\sim 40\%$  and remains reduced until the onset of the Holocene with values 315 halved with respect to present during the LGM. In the Gowan simulation, around 40 ka, the values increase to present-day levels and then decrease again to their LGM minimum. Overall, for  $M_2$ , the 317 simulations suggest that open ocean dissipation was enhanced for most of the last glacial period, but 318 that the exact magnitude is dependent on the sea level and ice sheet evolution. Less  $M_2$  energy was 319 lost in the shelf seas between 70 ka and the onset of the Holocene and more  $M_2$  energy dissipated in 320 the open ocean. Overall, the results from the ICE-6G simulations suggest that total dissipation was 321 larger than present for most of the glacial cycle, on average 25% (maximum  $\sim 40\%$ ), whereas for the 322 Gowan simulations which only extend back to 80 ka, the average increase was only 7% (maximum 323 28%). 324

For  $S_2$  (Figure 1c), in contrast, total dissipation is lower than present for most of the glacial cycle, except for the Holocene (for ICE-6G, the average decrease is 15%). This is driven by pronounced decreases in shelf sea dissipation which are strongest between 70 ka and the mid-deglacial. The magnitude of the decrease (up to 60% with respect to present) is slightly larger in the ICE-6G simulations due to the larger reduction in shelf area. Deep dissipation is slightly elevated (20–30%) in the ICE-6G simulation, but reduced with respect to present (apart from the LGM period) in the Gowan simulation.

For K<sub>1</sub> (Figure 1d), the overall pattern is similar to S<sub>2</sub>, with total dissipation being slightly reduced for most of the glacial cycle (for ICE-6G by  $\sim 10\%$ ) driven by lower levels of shelf sea dissipation. Again, open ocean dissipation in the ICE-6G simulation is enhanced between 70 ka and the onset of the Holocene and shelf sea dissipation remains reduced; whereas in the Gowan simulation, there are two distinct peaks situated around the sea level lowstands at 60 ka and the LGM. Between the lowstands, open ocean dissipation and shelf sea dissipation returns to presentday values.

For  $O_1$  (Figure 1e), the shelf sea dissipation signal is similar to  $S_2$  and  $K_1$ , but the decreases are compensated by increases in open ocean dissipation, such that total dissipation remains at values near present for most of the glacial cycle and an increase of up to 38% in total dissipation occurs
 during the peak of the LGM (cf. Sulzbach et al., 2023).

Overall, this means that, for the Gowan simulations, globally integrated dissipation, on average, was close to present-day levels between 80 ka and present, but for the ICE-6G simulations, total dissipation over the whole glacial cycle was on average ~15% larger than at present (Figure 1f). Deep dissipation in the ICE-6G runs was 57% greater on average (80% larger between 66 ka and 16 ka), and raised by 34% in the Gowan simulation. For shelf dissipation the mean reductions for ICE-6G and Gowan are 22 and 24%, respectively.

### 349 **3.2** Tidal evolution over the last 430 kyr

350

#### 3.2.1 Relationship between GMSL and tidal dissipation

For the ICE-6G simulations, open ocean dissipation is highly and significantly anti-correlated 35 with GMSL for all constituents (r < -0.88 for all constituents)—i.e., when sea level drops, open 352 ocean dissipation increases (see Fig 7). For shelf dissipation, similarly high magnitude signifi-353 cantly positive correlations emerge (r > 0.95) but the relationship is opposite, i.e., decreasing sea 354 levels reduce shelf dissipation. In contrast, total dissipation shows a weaker albeit still significant 355 relationship with GMSL, owing to the opposing relationships of open ocean and shelf dissipation. 356 For  $M_2$  and  $O_1$ , total dissipation is anti-correlated with GMSL (i.e., larger dissipation with lower 357 GMSL), whereas for  $S_2$  and  $K_1$  we observe the reverse. For the Gowan simulations, correlations 358 generally have the same sign as the correlations in the ICE-6G simulations, but with slightly reduced 359 magnitudes. Dissipation from the uniform sea level drop simulations shows similarly high negative 360 correlations between GMSL and open ocean dissipation as the ICE-6G simulations, and highly 361 positive correlations between GMSL and shelf dissipation. 362

Next, we evaluate the polynomial regression fit between GMSL and both shelf and open ocean 363 dissipation for each tidal constituent for all three sets of simulations, respectively (see Fig. 8). For 364 open ocean dissipation, all constituents apart from  $S_2$  show r > 0.95 between actual and estimated 365 dissipation values, and for shelf dissipation all regression models can explain more than 89% of the 366 variability. For the Gowan simulations, the regression fit is good for the diurnal constituents but 367 shows a slightly less good fit for the semi-diurnal constituents (r = 0.65 and r = 0.68 for M<sub>2</sub> and S<sub>2</sub>, 368 respectively). Regression-based total dissipation for each constituent (Figure S10 in Supplementary 369 Information S1), calculated as the sum of  $D_{deep}$  and  $D_{shelf}$  from the regression models for each 370 constituent, compares well with explicitly modelled dissipation (r > 0.90) for ICE-6G. A similarly good fit can be achieved with the uniform SL simulations. However, it is notable that for  $M_2$ , 372 the uniform SL simulations show only very small variations in dissipation (< 0.2 TW for GMSL 373 variations between -135 m to +20 m) in comparison to the ICE-6G simulations (> 1.1 TW for 374 GMSL variations between -130 m and +3 m), as increases in open ocean dissipation increases 375 are balanced by decreases in shelf dissipation. For the Gowan simulations, the regression models 376 are able to reproduce explicitly modelled values for all constituents well, apart from  $M_2$ , where the 377 regression model produces near constant total dissipation values and correlation between explicitly and regression total modelled dissipation is not significant. 379

Interestingly, pronounced differences emerge in the relationship between GMSL and  $M_2$  open 380 ocean dissipation (Fig. 8) across the sets of simulations. Whilst both ICE-6G and the uniform 381 SL simulations show a strongly negative relationship, the slope of the relationship is very different, 382 with the uniform SL simulations showing much smaller increases in dissipation than the ICE-6G simulations for the equivalent GMSL decrease. The Gowan simulations fall somewhat between the 384 two other estimates. These differences may reflect the strong control exerted by coastline positions on 385 the large  $M_2$  open ocean tides (e.g., Arbic et al., 2009; Green, 2010; Wilmes et al., 2019), especially 386 387 as differences in ice sheet extent and non-uniform sea level changes cause offsets in coastlines in comparison to the uniform sea level drop case. On the other hand, the non-uniform sea level changes 388 driven by GIA processes, which are especially pronounced close to ice sheets where we generally 389 see the largest tides, may also be contributing to these differences (see, e.g., Arbic et al., 2008). 390

#### 391 3.2.2 Dissipation estimates over the last 430 kyr

Following on, we now use the SL16 sea level curve and associated errors (Fig. 9a) to infer 392 dissipation back to 430 ka by applying all three regression models. The SL16 sea level curve shows 393 large variability in GMSL with sea level fluctuating between highstands of  $\sim -8$  m and +16 m during 394 the interglacials and sea level lowstands of 98–129 m during glacial maxima. On average, GMSL 395 was 54 m lower than present over the last 431 ka, according to SL16. The SL16 GMSL estimate 396 compares well with the ICE-6G derived GMSL for the last glacial cycle (r = 0.98, p = 0.00) 397 (Fig. 9a). Small offsets ( $\sim$ 15 m) can be seen during MIS 5, when climate entered into the last glacial phase. However, large differences emerge between SL16 and Gowan, which are especially pronounced during MIS 3 and late MIS 5. 400

Regression-inferred dissipation values for the last glacial cycle closely follow the explicitly 401 modelled values both for ICE-6G and Gowan (Fig. 9b-k). The largest offsets occur when the sea 402 level curves of ICE-6G and Gowan disagree with SL16, e.g., around the relative sea level highstands 403 at 100 ka and 80 ka for ICE-6G, and during MIS 3 for Gowan. The disparities are most pronounced 404 for  $M_2$  open ocean dissipation. Looking back over the last four glacial cycles, the ICE-6G regression 405 model suggests that  $M_2$  open ocean dissipation was strongly increased with respect to to present—on 406 average by 120%, i.e., more than doubled—apart from during interglacial phases. When GMSL 407 drops by more than 100 m,  $M_2$  open dissipation increases on average by a factor 2.6. For the Gowan 408 model, the mean  $M_2$  open ocean enhancements (36%) are less pronounced and and open ocean 409 dissipation during sea level lowstands increased by a factor 1.7. These Gowan change estimates are 410 very similar to those obtained using the uniform SL model (Figure S11 in Supplementary Information 411 S1). For shelf dissipation, all three models give very similar results, estimating the mean decrease in 412 dissipation to be around 20% over the last 430 kyr. Total  $M_2$  dissipation averaged over the last 430 413 kyr was between 3% (uniform SL), 5% (Gowan) and 37% (ICE-6G) greater than at present. For the 414 other constituents, average total dissipation changes are less pronounced, generally within  $\pm 25\%$ . 415 However, considerable fluctuations in the partitioning of tidal energy between the open ocean and the 416 shelf seas can be seen; with  $S_2$ ,  $K_1$  and  $O_1$  open ocean dissipation increasing during glacial periods 417 and decreasing towards present-day values during interglacials, respectively, and shelf dissipation behaving in an opposing manner. Taken together, these results suggest that globally integrated open 419 ocean dissipation was on average between 28 % (Gowan) and 85 % (ICE-6G) greater than at present, 420 with enhancements by a factor 1.5 to 2.2 during sea level lowstands and values similar to present 421 during highstands. However, it is worth noting that the exact magnitude of the estimated changes is 422 dependent on the specific sea level (and ice sheet) model adopted. 423

#### 424 **4 Discussion and Implications**

Our simulations are able to reproduce the findings of other studies for the LGM, deglacial and
Holocene (Arbic et al., 2004; Uehara et al., 2006; Griffiths & Peltier, 2008, 2009; Green, 2010;
Wilmes & Green, 2014; Wilmes et al., 2019, 2022), but also provide the first continuous estimates
of tides and tidal dissipation for the last 430 kyr, thus expanding our knowledge on tidal dynamics
prior to the LGM.

Our results show that, apart from interglacial phases, tides and tidal dissipation were different 430 from present during most of this extended time period, covering multiple glacial-interglacial cycles. 431 Changes were especially pronounced for the M<sub>2</sub> tidal constituent, which displays near-resonant 432 behaviour during the LGM (Arbic et al., 2004; Uehara et al., 2006; Griffiths & Peltier, 2008, 2009; 433 Green, 2010; Wilmes & Green, 2014; Wilmes et al., 2019). Our results show that the  $M_2$  tide in 434 the open ocean was strongly enhanced with respect to present for most of the last glacial cycles, 435 whereas for the constituents S<sub>2</sub>, K<sub>1</sub> and O<sub>1</sub> the relative enhancements were smaller and confined 436 to periods with the lowest sea levels. Peak open ocean amplitudes and dissipation occurred when 437 GMSL dropped below 70 m. Notably, for the LGM and for both the ICE-6G and Gowan simulations, peak dissipation values do not coincide with the lowest sea levels and greatest ice sheet extent but 439 reach their maximum 4-5 kyr after the lowstand as the ice begins to recede and sea level begins to 440 increase (at 22 ka for ICE-6G and 15 ka for Gowan). This is likely related to the location of grounded 441

ice margins and thus regional shelf area which in turn affects resonance properties of the glacial Atlantic (see experiments in Wilmes et al., 2019, which show that reduced Weddell Sea ice extent increases Atlantic dissipation by  $\sim 1$  TW).

However, there are considerable differences between the two reconstructions for the last glacial 445 cycle, and thus differences in the simulated tides. Whilst the two reconstructions show relatively 446 similar GMSL and ice sheet extent from 20 ka to present and agree during the two peak glaciations 447 (LGM and MIS 4), they strongly differ in ice sheet extent and GMSL during MIS 3 (55 ka to 25 448 ka) and late MIS 5 (80 ka to 65 ka) (see Fig. 2). The ICE-6G reconstruction shows an extensive 449 Laurentide ice sheet (LIS) from 80 ka and Fennoscandia is largely covered by ice from 70 ka onwards. 450 In contrast, the Gowan reconstruction has extensive ice sheets over North America and Fennoscandia 451 only during the glacial periods around 60 ka and the LGM. Ice sheet extent leading up to these periods 452 (80-65 ka and 55-30 ka) is much reduced in comparison to ICE-6G, and restricted to smaller more 453 localised ice caps (see Fig. 2). This leads to pronounced differences in local relative sea level of over 454 100 m and GMSL offsets of over 65 m between the two reconstructions during MIS 3 (see Figure 455 1a and Fig. 2). GMSL during MIS 3 is subject to high uncertainty with sea level estimates for this 456 periods remaining debated (e.g., Pico et al., 2016, 2017; Dalton et al., 2016, 2022). Because of a lack of direct markers of GMSL prior to the LGM, pre-LGM GMSL is often inferred from proxy records 458 such as  $\delta^{18}$ O from benthic foraminifera (e.g., Waelbroeck et al., 2002) or planktonic foraminifera 459 (Shakun et al., 2015), from spleotherms in the Red Sea (Grant et al., 2014; Rohling et al., 2009) 460 and compilations from multiple statistically analysed records (Spratt & Lisiecki, 2016). However, 461 because  $\delta^{18}$ O carries the imprint of both global ice-volume and ocean-temperature changes, the 462 records need to be corrected for temperature changes to obtain sea level curves (e.g., De Boer et 463 al., 2014). Large offsets in MIS 3 GMSL of 30-60 m have been found between ice sheet based 464 reconstructions (e.g., Pico et al., 2016, 2017; Dalton et al., 2016; Gowan et al., 2021; Dalton et 465 al., 2022) and records inferred from marine  $\delta^{18}$ O (e.g., Waelbroeck et al., 2002; Spratt & Lisiecki, 466 2016), with ice sheet reconstructions suggesting that (i) the Laurentide Ice Sheet during MIS 3 was 467 much less extensive than previously thought and (ii) GMSL was  $\sim$ 30–50 m lower than at present 468 (Pico et al., 2016, 2017; Dalton et al., 2016, 2022). This implies that, for MIS 3, the Gowan sea level 469 and ice sheet reconstruction may be more appropriate as it captures the relative sea level highstand 470 and reduced ice sheet extent. 471

By comparing the ICE-6G and Gowan tide simulations for the LGM, it becomes apparent how 472 sensitive the tidal dynamics are to relative small changes in bathymetry, i.e., changes in both sea level 473 and ocean basin shape (e.g., through ice sheet extent changes) when in a near-resonant state. At 20 474 ka,  $M_2$  open ocean dissipation differs by 0.5 TW between the ICE-6G and Gowan simulation, despite 475 a very similar GMSL drop of 116 m and 120 m, respectively. Experiments in Wilmes et al. (2019) 476 showed that LGM M<sub>2</sub> tidal dynamics were remarkably insensitive to offsets in GMSL ( $\pm 10$  m) but that changes in the location of the land-ocean boundaries could have dramatic effects (e.g., altering 478 ice sheet extent in the Weddell Sea led to a > 1 TW change in M<sub>2</sub> dissipation). This is also shown 479 when comparing the uniform SL change simulations with those that have realistic glacial land ocean 480 boundaries (1.8 TW versus 2.5 and 2.0 TW for ICE-6G and Gowan, respectively). Comparing the 481 ICE-6G and Gowan 20 ka bathymetries (see Fig. 2) shows that the Gowan reconstruction has more 482 ice around the margins of the Laurentide and the Greenland Ice Sheets but also around the margins 483 of Antarctica. Furthermore, whilst GMSL is very similar between the two reconstructions, regional sea levels in the Labrador Sea differ by over 50 m (Fig. 2) which may also influence tidal dynamics 485 and contribute to the dissipation offsets. It is also worth noting that the LGM sea level lowstand 486 does not occur at the same time in the two reconstructions: in ICE-6G, the lowest sea levels occur 487 around 26 ka, whereas for Gowan they occur at 20 ka, and the ICE-6G lowstand is 10 m lower than 488 in Gowan (-130 m versus -120 m). 489

This work and that of Wilmes et al. (2019) demonstrate that small changes in land-ocean boundaries, e.g., through changes in ice sheet extent, could dramatically alter open ocean tidal energy dissipation on the order of  $\pm 1$  TW. In experiments using the intermediate climate model UVic (Schmittner et al., 2015; Wilmes et al., 2019, 2021), changes of this order of magnitude correspond to a ~2–4 Sv change in AMOC strength. These strong sensitivities in glacial tides could

cause rapid changes in the amount of tidal energy available for ocean mixing which, in turn, could 495 impact ocean circulation and thus climate. Feedbacks between tides, ice sheet extent, ocean mixing, 496 and climate may be a modulating factor for glacial climate, especially for period such as Heinrich 497 events, where ice loss from Hudson Strait has been postulated to have affected tidal dynamics and vice versa (e.g., Arbic et al., 2004; Velay-Vitow et al., 2020). For instance, in the Gowan simulations, 499 between 60 ka and 25 ka, fluctuations in  $M_2$  open ocean dissipation can be seen which appear 500 related to Hudson Bay & Strait dissipation, thus suggesting that Hudson Bay & Strait ice cover may 501 affect (North) Atlantic tidal energy availability. Furthermore, the large tides along the ice sheet 502 margins—both for the Antarctic and Northern Hemisphere ice sheets—may have contributed to the 503 rapid ice sheet retreat linked with the deglacial increase in sea level (Gomez et al., 2020; Batchelor 504 et al., 2023), for example, through ice shelf tidal flexure and crevasse formation (e.g., Olinger et al., 505 2019). 506

Because the sea level simulations used for the last glacial cycle include no ice history prior to 507 122 ka (ICE-6G) and 80 ka (Gowan), the oldest timeslices (i.e., LIG for ICE-6G and late MIS 5 for 508 Gowan) are likely associated with the largest errors, as each time step is affected by previous ice sheet 509 and sea level history. For example, because no ice history was present for the penultimate glacial 510 maximum, the exact sea levels and thus tides during the LIG, are likely subject to larger errors. It 511 is also possible that, due to ice history 'memory' from previous glaciations, tidal dynamics moving 512 into a glaciation may behave differently than during the deglacial phase. However, at present, due to 513 a lack of global high resolution ice sheet reconstructions beyond the LIG, it is not possible to evaluate 514 this point. Another source of uncertainty is the relatively low temporal resolution of the ice sheet 515 histories (2.5 kyr for Gowan throughout the last glacial cycle and 2 kyr prior to 32 ka for ICE-6G, 516  $\leq$  1 kyr after 32 ka) and for ICE-6G, the horizontal resolution of the ice sheets (1/2° × 1/2°). It is 517 conceivable that the adopted histories lack small scale ice sheet and bathymetry fluctuations, which, 518 during tidally resonant states, could have led to large short-term regional and supraregional shifts in 519 tidal dynamics. 520

Simulations of climate during the last glacial cycle or parts thereof generally neglect transient 521 changes in tidal energy input to ocean mixing and hold mixing rates constant at the present-day or glacial state (e.g., Menviel et al., 2017; Pöppelmeier et al., 2023). However, this work emphasises 523 that tidal dissipation strongly varied between glacial and interglacial phases and that large short-524 term variations may have occurred during transitions into and out of glacial phases. The explicitly 525 simulated glacial cycle tidal dissipation timeslices can be used to adjust vertical mixing in transient 526 climate model simulations. These also account for regional responses in dissipation which may 527 differ from the global mean (e.g., during MIS 3 dissipation dissipation enhancements in the Southern 528 Hemisphere are greater than in the Northern Hemisphere which stands in contrast to full glacial 529 conditions during the LGM). 530

#### 531 5 Conclusion

In this study, we have simulated tidal amplitudes and tidal dissipation for the last glacial cycle 532 encompassing the LIG to present. In addition, we have estimated tidal dissipation for the last four 533 glacial cycles spanning back to 430 ka. Our findings suggest that tides likely differed from present 534 during most glacial phases, with maximum tidal amplitudes and open ocean dissipation occurring 535 during sea level lowstands due to the resonant properties of the glacial ocean with regards to the 536 semi-diurnal tidal forcing. Additionally, we discovered that during glacial phases, semi-diurnal tides, 537 particularly  $M_2$ , were highly sensitive to sea level and ice sheet extent, which remain uncertain prior 600 to the LGM. These results are of wider relevance as they indicate that the availability of tidal energy during the late Pleistocene strongly differed from the present, potentially impacting ocean mixing, 540 and thus ocean circulation and climate in the past. 541

### 542 6 Open Research

The data used in this study are available using the following links: Gowan et al. (2021) ice sheet and sea level reconstruction https://doi.pangaea.de/10.1594/PANGAEA.905800, ICE-6G ice sheet reconstruction https://www.atmosp.physics.utoronto.ca/~peltier/data.php,
 TPXO9 tidal atlas https://www.tpxo.net/global/tpxo9-atlas and RTopo-2 https://doi
 .pangaea.de/10.1594/PANGAEA.856844. The tidal simulations and relevant grid files together
 with the Matlab routines necessary to read the data are available at the repository Zenodo via
 doi:10.5281/zenodo.8147509 (https://zenodo.org/record/8147509) with open access (Wilmes
 et al., 2023).

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**Figure 1.** (a) Global mean sea level (black line) and shelf area (in % of global ocean area) for ICE-6G (solid lines) and Gowan (2021) (dotted lines). Marine isotope stages 1 to 5 are indicated by shading and labels above the timeline. For MIS 5, letters indicate substages. Interglacial periods are highlighted by yellow shading. Italic labels show climate periods: LIG = Last Interglacial; LGM = Last Glacial Maximum. (b–e) Dissipation for the constituents  $M_2$ ,  $S_2$ ,  $K_1$  and  $O_1$ , respectively. Open ocean (deep) dissipation is plotted in blue, shelf sea dissipation in red and globally integrated dissipation in black. ICE-6G values are plotted with solid lines and Gowan (2021) values with dashed lines. (f) same as (b–e) but for the sum of the constituents. The dashed thin straight lines give present-day values as a reference.























**Figure 7.** Correlation coefficients between GMSL and dissipation by tidal constituent from runs with ICE-6G bathymetry (circular markers), with uniform sea level changes (squares) and the Gowan bathymetries (crosses). Blue markers are used for open ocean dissipation, red ones for shelf sea dissipation and black ones for globally integrated dissipation.



**Figure 8.** Open ocean dissipation (blue, left column) and shelf sea dissipation (red, right column) plotted against GMSL for the constituents  $M_2$ ,  $S_2$ ,  $K_1$  and  $O_1$  (top to bottom). Circular makers show dissipation from the ICE-6G simulations, Gowan simulations are shown with crosses, and squares plot dissipation from the uniform SL runs ( $M_2$  and  $K_1$  only). Regression lines are plotted in black (solid for ICE-6G and dashed for uniform SL runs). The correlation coefficients *r* between the explicitly modelled and regression-estimated dissipation values are printed in each panel for ICE-6G ( $r_I$ ), the uniform SL simulations ( $r_U$ ), and the Gowan simulations ( $r_G$ ). Stars indicate that the correlation is significant at the 99% confidence level.



**Figure 9.** (a) Global mean sea level from the Spratt and Lisiecki (2016) reconstruction (blue line) with shading indicating upper and lower uncertainties, from the ICE-6G bathymetry (black solid line) and from Gowan et al. (2021) (black dotted line). Light blue and light yellow shading with corresponding numbering indicates Marine Isotope Stages. (b, d, f, h, j) Regression estimated dissipation for the constituents  $M_2$ ,  $S_2$ ,  $K_1$  and  $O_1$ , respectively, and the sum thereof. Open ocean (deep) dissipation is plotted in blue, shelf sea dissipation in red and globally integrated dissipation in black using the ICE-6G regression models. Shading indicates one standard deviation uncertainty of the dissipation estimates. Explicitly modelled dissipation values for the last glacial cycle are overlain in with dashed lines. Light blue and light yellow shading with corresponding numbering indicates Marine Isotope Stages. (c, e, g, i, j, k) same as (b, d, f, h, j) but for the Gowan et al. (2021) regression model.

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